

# An introduction to Discrete de Rham (DDR) methods

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1 Three model problems and their well-posedness

2 Discrete de Rham (DDR) complexes

3 Application to magnetostatics

# Setting I

- Let  $\Omega \subset \mathbb{R}^3$  be a connected polyhedral domain with **Betti numbers**  $b_i$
- We have  $b_0 = 1$  (number of connected components) and  $b_3 = 0$
- $b_1$  accounts for the number of **tunnels** crossing  $\Omega$



$$(b_0, b_1, b_2, b_3) = (1, \mathbf{1}, 0, 0)$$

- $b_2$  is the number of **voids** encapsulated by  $\Omega$



$$(b_0, b_1, b_2, b_3) = (1, 0, \mathbf{1}, 0)$$

# Setting II

- We consider PDE models that hinge on the **vector calculus operators**:

$$\mathbf{grad} \, q = \begin{pmatrix} \partial_1 q \\ \partial_2 q \\ \partial_3 q \end{pmatrix}, \quad \mathbf{curl} \, \mathbf{v} = \begin{pmatrix} \partial_2 v_3 - \partial_3 v_2 \\ \partial_3 v_1 - \partial_1 v_3 \\ \partial_1 v_2 - \partial_2 v_1 \end{pmatrix}, \quad \operatorname{div} \mathbf{w} = \partial_1 w_1 + \partial_2 w_2 + \partial_3 w_3$$

for smooth enough functions

$$q : \Omega \rightarrow \mathbb{R}, \quad \mathbf{v} : \Omega \rightarrow \mathbb{R}^3, \quad \mathbf{w} : \Omega \rightarrow \mathbb{R}^3$$

- The corresponding  $L^2$ -graph (domain) spaces are

$$\begin{aligned} H^1(\Omega) &:= \{q \in L^2(\Omega) : \mathbf{grad} \, q \in L^2(\Omega) := L^2(\Omega)^3\}, \\ \mathbf{H}(\mathbf{curl}; \Omega) &:= \{\mathbf{v} \in L^2(\Omega) : \mathbf{curl} \, \mathbf{v} \in L^2(\Omega)\}, \\ \mathbf{H}(\operatorname{div}; \Omega) &:= \{\mathbf{w} \in L^2(\Omega) : \operatorname{div} \mathbf{w} \in L^2(\Omega)\} \end{aligned}$$

- Assume for the moment that  $\Omega$  has **trivial topology** (i.e.,  $b_1 = b_2 = 0$ )

# Three model problems

## The Stokes problem in curl-curl formulation

- Given a real number  $\nu > 0$  and  $\mathbf{f} \in \mathbf{L}^2(\Omega)$ , the Stokes problem reads:  
Find the **velocity**  $\mathbf{u} : \Omega \rightarrow \mathbb{R}^3$  and **pressure**  $p : \Omega \rightarrow \mathbb{R}$  s.t.

$$\begin{aligned} \overbrace{\nu(\operatorname{curl} \operatorname{curl} \mathbf{u} - \operatorname{grad} \operatorname{div} \mathbf{u})}^{-\nu \Delta \mathbf{u}} + \operatorname{grad} p &= \mathbf{f} && \text{in } \Omega, && \text{(momentum conservation)} \\ \operatorname{div} \mathbf{u} &= 0 && \text{in } \Omega, && \text{(mass conservation)} \\ \operatorname{curl} \mathbf{u} \times \mathbf{n} = \mathbf{0} \text{ and } \mathbf{u} \cdot \mathbf{n} &= 0 && \text{on } \partial\Omega, && \text{(boundary conditions)} \\ \int_{\Omega} p &= 0 \end{aligned}$$

- Weak formulation:** Find  $(\mathbf{u}, p) \in \mathbf{H}(\operatorname{curl}; \Omega) \times H^1(\Omega)$  s.t.  $\int_{\Omega} p = 0$  and

$$\begin{aligned} \int_{\Omega} \nu \operatorname{curl} \mathbf{u} \cdot \operatorname{curl} \mathbf{v} + \int_{\Omega} \operatorname{grad} p \cdot \mathbf{v} &= \int_{\Omega} \mathbf{f} \cdot \mathbf{v} && \forall \mathbf{v} \in \mathbf{H}(\operatorname{curl}; \Omega), \\ - \int_{\Omega} \mathbf{u} \cdot \operatorname{grad} q &= 0 && \forall q \in H^1(\Omega) \end{aligned}$$

# Three model problems

## The magnetostatics problem

- For  $\mu > 0$  and  $\mathbf{J} \in \mathbf{curl} \mathbf{H}(\mathbf{curl}; \Omega)$ , the magnetostatics problem reads:  
Find the **magnetic field**  $\mathbf{H} : \Omega \rightarrow \mathbb{R}^3$  and **vector potential**  $\mathbf{A} : \Omega \rightarrow \mathbb{R}^3$  s.t.

$$\begin{aligned}\mu \mathbf{H} - \mathbf{curl} \mathbf{A} &= \mathbf{0} && \text{in } \Omega, && \text{(vector potential)} \\ \mathbf{curl} \mathbf{H} &= \mathbf{J} && \text{in } \Omega, && \text{(Ampère's law)} \\ \operatorname{div} \mathbf{A} &= 0 && \text{in } \Omega, && \text{(Coulomb's gauge)} \\ \mathbf{A} \times \mathbf{n} &= \mathbf{0} && \text{on } \partial\Omega && \text{(boundary condition)}\end{aligned}$$

- Weak formulation:** Find  $(\mathbf{H}, \mathbf{A}) \in \mathbf{H}(\mathbf{curl}; \Omega) \times \mathbf{H}(\operatorname{div}; \Omega)$  s.t.

$$\begin{aligned}\int_{\Omega} \mu \mathbf{H} \cdot \boldsymbol{\tau} - \int_{\Omega} \mathbf{A} \cdot \mathbf{curl} \boldsymbol{\tau} &= 0 && \forall \boldsymbol{\tau} \in \mathbf{H}(\mathbf{curl}; \Omega), \\ \int_{\Omega} \mathbf{curl} \mathbf{H} \cdot \mathbf{v} + \int_{\Omega} \operatorname{div} \mathbf{A} \operatorname{div} \mathbf{v} &= \int_{\Omega} \mathbf{J} \cdot \mathbf{v} && \forall \mathbf{v} \in \mathbf{H}(\operatorname{div}; \Omega)\end{aligned}$$

# Three model problems

## The Darcy problem in velocity-pressure formulation

- Given  $\kappa > 0$  and  $f \in L^2(\Omega)$ , the Darcy problem reads:

Find the **velocity**  $\mathbf{u} : \Omega \rightarrow \mathbb{R}^3$  and **pressure**  $p : \Omega \rightarrow \mathbb{R}$  s.t.

$$\begin{aligned}\kappa^{-1} \mathbf{u} - \mathbf{grad} \, p &= 0 && \text{in } \Omega, && (\text{Darcy's law}) \\ -\operatorname{div} \mathbf{u} &= f && \text{in } \Omega, && (\text{mass conservation}) \\ p &= 0 && \text{on } \partial\Omega && (\text{boundary condition})\end{aligned}$$

- Weak formulation:** Find  $(\mathbf{u}, p) \in \mathbf{H}(\operatorname{div}; \Omega) \times L^2(\Omega)$  s.t.

$$\begin{aligned}\int_{\Omega} \kappa^{-1} \mathbf{u} \cdot \mathbf{v} + \int_{\Omega} p \operatorname{div} \mathbf{v} &= 0 && \forall \mathbf{v} \in \mathbf{H}(\operatorname{div}; \Omega), \\ -\int_{\Omega} \operatorname{div} \mathbf{u} \, q &= \int_{\Omega} f \, q && \forall q \in L^2(\Omega)\end{aligned}$$



# A unified view

- The above problems are **mixed formulations** involving two fields
- They can be recast into the abstract setting: Find  $(\sigma, u) \in \Sigma \times U$  s.t.

$$\begin{aligned}a(\sigma, \tau) + b(\tau, u) &= f(\tau) \quad \forall \tau \in \Sigma, \\ -b(\sigma, v) + c(u, v) &= g(v) \quad \forall v \in U,\end{aligned}$$

or, equivalently, in variational formulation,

$$\mathcal{A}((\sigma, u), (\tau, v)) = f(\tau) + g(v) \quad \forall (\tau, v) \in \Sigma \times U$$

with

$$\mathcal{A}((\sigma, u), (\tau, v)) := a(\sigma, \tau) + b(\tau, u) - b(\sigma, v) + c(u, v) = f(\tau) + g(v)$$

- Well-posedness holds under an **inf-sup condition on  $\mathcal{A}$**

# A unified tool for well-posedness: The de Rham complex

$$\{0\} \longrightarrow H^1(\Omega) \xrightarrow{\text{grad}} \mathbf{H}(\text{curl}; \Omega) \xrightarrow{\text{curl}} \mathbf{H}(\text{div}; \Omega) \xrightarrow{\text{div}} L^2(\Omega) \longrightarrow \{0\}$$

- Key properties, possibly depending on the topology of  $\Omega$ :

$$\text{Im } \mathbf{grad} \subset \text{Ker } \mathbf{curl},$$

$$\text{Im } \mathbf{curl} \subset \text{Ker } \text{div},$$

$$\Omega \subset \mathbb{R}^3 \ (b_3 = 0) \implies \text{Im } \text{div} = L^2(\Omega) \quad (\text{Darcy, magnetostatics})$$

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- Key properties, possibly depending on the topology of  $\Omega$ :

no tunnels crossing  $\Omega$  ( $b_1 = 0$ )  $\implies \text{Im grad} = \text{Ker curl}$  (Stokes)

no voids contained in  $\Omega$  ( $b_2 = 0$ )  $\implies \text{Im curl} = \text{Ker div}$  (magnetostatics)

$\Omega \subset \mathbb{R}^3$  ( $b_3 = 0$ )  $\implies \text{Im div} = L^2(\Omega)$  (Darcy, magnetostatics)

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no voids contained in  $\Omega$  ( $b_2 = 0$ )  $\implies$  **Im curl = Ker div** (magnetostatics)

$\Omega \subset \mathbb{R}^3$  ( $b_3 = 0$ )  $\implies$  **Im div =  $L^2(\Omega)$**  (Darcy, magnetostatics)

- When  $b_1 \neq 0$  or  $b_2 \neq 0$ , **de Rham's cohomology** characterizes

$$\text{Ker curl} / \text{Im grad} \quad \text{and} \quad \text{Ker div} / \text{Im curl}$$

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$$\text{Ker } \text{curl} / \text{Im } \text{grad} \quad \text{and} \quad \text{Ker } \text{div} / \text{Im } \text{curl}$$

- Emulating these algebraic properties is key for stable discretizations

# Generalization through differential forms

- The de Rham complex generalizes to **domains of  $\mathbb{R}^n$**  or smooth manifolds
- Denoting by  $d$  the **exterior derivative** and by  $H\Lambda(\Omega)$  its  $L^2$ -domain,

$$H\Lambda^0(\Omega) \xrightarrow{d^0} \cdots \xrightarrow{d^{k-1}} H\Lambda^k(\Omega) \xrightarrow{d^k} \cdots \xrightarrow{d^{n-1}} H\Lambda^n(\Omega) \longrightarrow \{0\}$$

- For  $n = 3$ , the vector calculus version is recovered through **vector proxies**

$$\begin{array}{ccccccc} H\Lambda^0(\Omega) & \xrightarrow{d^0} & H\Lambda^1(\Omega) & \xrightarrow{d^1} & H\Lambda^2(\Omega) & \xrightarrow{d^2} & H\Lambda^3(\Omega) \longrightarrow \{0\} \\ \updownarrow \cong & & \updownarrow \cong & & \updownarrow \cong & & \updownarrow \cong \\ H^1(\Omega) & \xrightarrow{\text{grad}} & \mathbf{H}(\text{curl}; \Omega) & \xrightarrow{\text{curl}} & \mathbf{H}(\text{div}; \Omega) & \xrightarrow{\text{div}} & L^2(\Omega) \longrightarrow \{0\} \end{array}$$

# The (trimmed) Finite Element way

## Local spaces

- Let  $T \subset \mathbb{R}^3$  be a polyhedron and set, for any  $k \geq -1$ ,

$$\mathcal{P}^k(T) := \{\text{restrictions of 3-variate polynomials of degree } \leq k \text{ to } T\}$$

- Fix  $k \geq 0$  and write, denoting by  $\mathbf{x}_T$  a point inside  $T$ ,

$$\begin{aligned}\mathcal{P}^k(T)^3 &= \mathbf{grad} \mathcal{P}^{k+1}(T) \oplus (\mathbf{x} - \mathbf{x}_T) \times \mathcal{P}^{k-1}(T)^3 =: \mathcal{G}^k(T) \oplus \mathcal{G}^{c,k}(T) \\ &= \mathbf{curl} \mathcal{P}^{k+1}(T)^3 \oplus (\mathbf{x} - \mathbf{x}_T) \mathcal{P}^{k-1}(T) =: \mathcal{R}^k(T) \oplus \mathcal{R}^{c,k}(T)\end{aligned}$$

- Define the **trimmed spaces** that sit between  $\mathcal{P}^k(T)^3$  and  $\mathcal{P}^{k+1}(T)^3$ :

$$\begin{aligned}\mathcal{N}^{k+1}(T) &:= \mathcal{G}^k(T) \oplus \mathcal{G}^{c,k+1}(T) \\ \mathcal{RT}^{k+1}(T) &:= \mathcal{R}^k(T) \oplus \mathcal{R}^{c,k+1}(T)\end{aligned}$$

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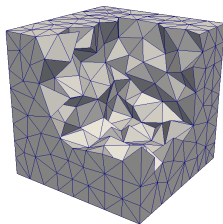
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- $\mathcal{P}^{-,k} \Lambda^r(f)$  generalizes to  $r$ -forms on  $d$ -faces  $f$  through **Koszul complements**



# The (trimmed) Finite Element way

Global complex



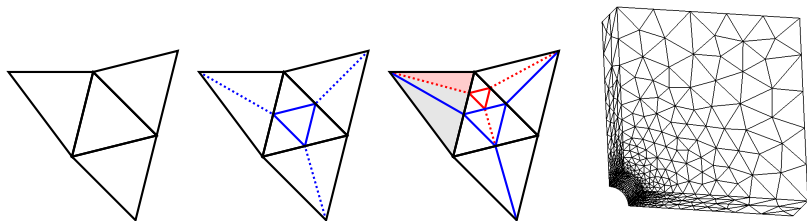
- Let  $\mathcal{T}_h$  be a **conforming tetrahedral mesh** of  $\Omega$  and let  $k \geq 0$
- Local spaces can be **glued together** to form a **global FE complex**:

$$\begin{array}{ccccccccc} \mathcal{P}_c^{k+1}(\mathcal{T}_h) & \xrightarrow{\text{grad}} & \mathcal{N}^{k+1}(\mathcal{T}_h) & \xrightarrow{\text{curl}} & \mathcal{RT}^{k+1}(\mathcal{T}_h) & \xrightarrow{\text{div}} & \mathcal{P}^k(\mathcal{T}_h) & \longrightarrow & \{0\} \\ \downarrow & & \downarrow & & \downarrow & & \downarrow & & \\ H^1(\Omega) & \xrightarrow{\text{grad}} & H(\text{curl}; \Omega) & \xrightarrow{\text{curl}} & H(\text{div}; \Omega) & \xrightarrow{\text{div}} & L^2(\Omega) & \longrightarrow & \{0\} \end{array}$$

- **The gluing only works on conforming meshes (simplicial complexes)!**

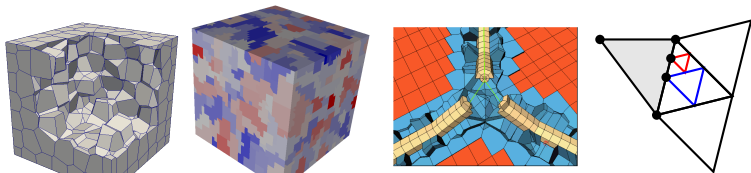
# The Finite Element way

## Shortcomings



- Approach limited to **conforming meshes** with **standard elements**
  - ⇒ local refinement requires to **trade mesh size for mesh quality**
  - ⇒ complex geometries may require a **large number of elements**
  - ⇒ the element shape cannot be **adapted to the solution**
- Need for (global) basis functions
  - ⇒ significant increase of DOFs on hexahedral elements

# The discrete de Rham (DDR) approach I

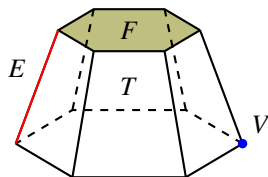


- **Key idea:** replace both spaces and operators by discrete counterparts:

$$\underline{X}_{\text{grad},h}^k \xrightarrow{\underline{G}_h^k} \underline{X}_{\text{curl},h}^k \xrightarrow{\underline{C}_h^k} \underline{X}_{\text{div},h}^k \xrightarrow{D_h^k} \mathcal{P}^k(\mathcal{T}_h) \longrightarrow \{0\}$$

- Support of **polyhedral meshes** (CW complexes) and **high-order**
- Several strategies to **reduce the number of unknowns** on general shapes
- Natural generalization to the **de Rham complex of differential forms**

# The discrete de Rham (DDR) approach II



- DDR spaces are spanned by **vectors of polynomials**
- Polynomial components enable **consistent reconstructions** of
  - Vector calculus operators
  - The corresponding scalar or vector potentials
- These reconstructions emulate **integration by parts (Stokes) formulas**

# References for this presentation

- Vector FE spaces [Raviart and Thomas, 1977, Nédélec, 1980]
- FEEC [Arnold, Falk, Winther, 2006–pres.]
- Introduction of DDR [DP, Droniou, Rapetti, 2020]
- DDR with Koszul complements [DP and Droniou, 2023a]
- Algebraic properties (general topology) [DP, Droniou, Pitassi, 2023]
- Bridges with VEM [Beirão da Veiga, Dassi, DP, Droniou, 2022]
- Polytopal Exterior Calculus (PEC) [Bonaldi, DP, Droniou, Hu, 2023]
- C++ open-source implementation available in **HArDCore3D**

# Outline

1 Three model problems and their well-posedness

2 Discrete de Rham (DDR) complexes

3 Application to magnetostatics

# The two-dimensional case

## Continuous exact complex

- With  **$F$  mesh face** let, for  $q : F \rightarrow \mathbb{R}$  and  $\mathbf{v} : F \rightarrow \mathbb{R}^2$  smooth enough,

$$\mathbf{rot}_F q := (\mathbf{grad}_F q)^\perp \quad \mathbf{rot}_F \mathbf{v} := \mathbf{div}_F(\mathbf{v}^\perp)$$

- We start by deriving a discrete counterpart of the **2D de Rham complex**:

$$H^1(F) \xrightarrow{\mathbf{grad}_F} \mathbf{H}(\mathbf{rot}; F) \xrightarrow{\mathbf{rot}_F} L^2(F) \longrightarrow \{0\}$$

- We will need the following decomposition of  $\mathcal{P}^k(F)^2$ :

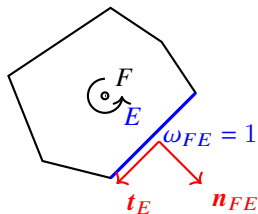
$$\mathcal{P}^k(F)^2 = \mathbf{rot}_F \mathcal{P}^{k+1}(F) \oplus (\mathbf{x} - \mathbf{x}_F) \mathcal{P}^{k-1}(F) =: \mathcal{R}^k(F) \oplus \mathcal{R}^{c,k}(F),$$

and recall the 2D Raviart–Thomas space

$$\mathcal{RT}^{k+1}(F) := \mathcal{R}^k(F) \oplus \mathcal{R}^{c,k+1}(F)$$

# The two-dimensional case

A key remark



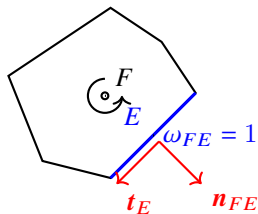
- Let  $q \in \mathcal{P}^{k+1}(F)$ . For any  $\mathbf{v} \in \mathcal{P}^k(F)^2$ , we have

$$\int_F \mathbf{grad}_F q \cdot \mathbf{v} = - \int_F q \operatorname{div}_F \mathbf{v} + \sum_{E \in \mathcal{E}_F} \omega_{FE} \int_E q|_{\partial F} (\mathbf{v} \cdot \mathbf{n}_{FE})$$



# The two-dimensional case

A key remark

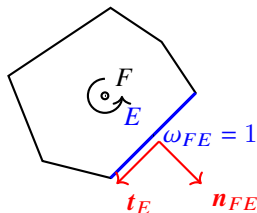


- Let  $q \in \mathcal{P}^{k+1}(F)$ . For any  $\mathbf{v} \in \mathcal{P}^k(F)^2$ , we have

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# The two-dimensional case

A key remark



- Let  $q \in \mathcal{P}^{k+1}(F)$ . For any  $\mathbf{v} \in \mathcal{P}^k(F)^2$ , we have

$$\int_F \mathbf{grad}_F q \cdot \mathbf{v} = - \int_F \underbrace{\pi_{\mathcal{P},F}^{k-1} q}_{\in \mathcal{P}^{k-1}(F)} \operatorname{div}_F \mathbf{v} + \sum_{E \in \mathcal{E}_F} \omega_{FE} \int_E q|_{\partial F} (\mathbf{v} \cdot \mathbf{n}_{FE})$$

- Hence,  $\mathbf{grad}_F q$  can be computed given  $\pi_{\mathcal{P},F}^{k-1} q$  and  $q|_{\partial F}$

# The two-dimensional case

## Discrete $H^1(F)$ space

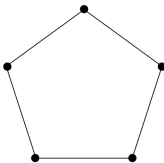
- Based on this remark, we take as discrete counterpart of  $H^1(F)$

$$\underline{X}_{\text{grad},F}^k := \left\{ \underline{q}_F = (q_F, q_{\partial F}) : q_F \in \mathcal{P}^{k-1}(F) \text{ and } q_{\partial F} \in \mathcal{P}_c^{k+1}(\mathcal{E}_F) \right\}$$

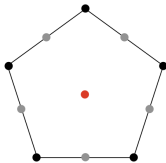
- Let  $\underline{I}_{\text{grad},F}^k : C^0(\overline{F}) \rightarrow \underline{X}_{\text{grad},F}^k$  be s.t.,  $\forall q \in C^0(\overline{F})$ ,

$$\underline{I}_{\text{grad},F}^k q := (\pi_{\mathcal{P},F}^{k-1} q, q_{\partial F}) \text{ with}$$

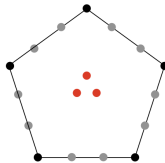
$$\pi_{\mathcal{P},E}^{k-1}(q_{\partial F})|_E = \pi_{\mathcal{P},E}^{k-1} q|_E \quad \forall E \in \mathcal{E}_F \text{ and } q_{\partial F}(\mathbf{x}_V) = q(\mathbf{x}_V) \quad \forall V \in \mathcal{V}_F$$



$k = 0$



$k = 1$



$k = 2$

# The two-dimensional case

Reconstructions in  $\underline{X}_{\text{grad},F}^k$

- For all  $E \in \mathcal{E}_F$ , the **edge gradient**  $G_E^k : \underline{X}_{\text{grad},F}^k \rightarrow \mathcal{P}^k(E)$  is s.t.

$$G_E^k \underline{q}_F := (q_{\partial F})'|_E$$

- The **face gradient**  $\mathbf{G}_F^k : \underline{X}_{\text{grad},F}^k \rightarrow \mathcal{P}^k(F)^2$  is s.t.,  $\forall \mathbf{v} \in \mathcal{P}^k(F)^2$ ,

$$\int_F \mathbf{G}_F^k \underline{q}_F \cdot \mathbf{v} = - \int_F q_F \operatorname{div}_F \mathbf{v} + \sum_{E \in \mathcal{E}_F} \omega_{FE} \int_E q_{\partial F} (\mathbf{v} \cdot \mathbf{n}_{FE})$$

- The **scalar trace**  $\gamma_F^{k+1} : \underline{X}_{\text{grad},F}^k \rightarrow \mathcal{P}^{k+1}(F)$  is s.t., for all  $\mathbf{v} \in \mathcal{R}^{c,k+2}(F)$ ,

$$\int_F \gamma_F^{k+1} \underline{q}_F \operatorname{div}_F \mathbf{v} = - \int_F \mathbf{G}_F^k \underline{q}_F \cdot \mathbf{v} + \sum_{E \in \mathcal{E}_F} \omega_{FE} \int_F q_{\partial F} (\mathbf{v} \cdot \mathbf{n}_{FE})$$

- By construction, we have **polynomial consistency**: For all  $q \in \mathcal{P}^{k+1}(F)$ ,

$$\mathbf{G}_F^k (\underline{I}_{\text{grad},F}^k q) = \operatorname{grad}_F q \text{ and } \gamma_F^{k+1} (\underline{I}_{\text{grad},F}^k q) = q$$

# The two-dimensional case

Discrete  $\mathbf{H}(\text{rot}; F)$  space

- We start from:  $\forall \mathbf{v} \in \mathcal{N}^{k+1}(F) := \mathcal{RT}^{k+1}(F)^\perp, \forall q \in \mathcal{P}^k(F),$

$$\int_F \text{rot}_F \mathbf{v} \, q = \int_F \mathbf{v} \cdot \underbrace{\text{rot}_F q}_{\in \mathcal{RT}^k(F)} - \sum_{E \in \mathcal{E}_F} \omega_{FE} \int_E (\mathbf{v} \cdot \mathbf{t}_E) q|_E$$

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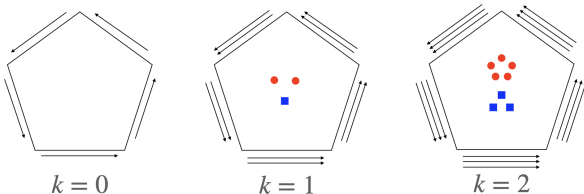
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- This leads to the following discrete counterpart of  $\mathbf{H}(\text{rot}; F)$ :

$$\underline{\mathbf{X}}_{\text{curl}, F}^k := \left\{ \underline{\mathbf{v}}_F = (\mathbf{v}_F, (\mathbf{v}_E)_{E \in \mathcal{E}_F}) : \mathbf{v}_F \in \mathcal{RT}^k(F) \text{ and } \mathbf{v}_E \in \mathcal{P}^k(E) \forall E \in \mathcal{E}_F \right\}$$

- $\underline{\mathbf{I}}_{\text{rot}, F}^k : H^1(F)^2 \rightarrow \underline{\mathbf{X}}_{\text{curl}, F}^k$  is obtained collecting  $L^2$ -orthogonal projections



# The two-dimensional case

Reconstructions in  $\underline{X}_{\text{curl},F}^k$

- The **face curl operator**  $C_F^k : \underline{X}_{\text{curl},F}^k \rightarrow \mathcal{P}^k(F)$  is s.t.,

$$\int_F \mathbf{C}_{F \underline{\mathbf{v}}_F}^k q = \int_F \mathbf{v}_F \cdot \mathbf{rot}_F q - \sum_{E \in \mathcal{E}_F} \omega_{FE} \int_E \mathbf{v}_E q \quad \forall q \in \mathcal{P}^k(F)$$

- The **tangent trace**  $\gamma_{t,F}^k : \underline{X}_{\text{curl},F}^k \rightarrow \mathcal{P}^k(F)^2$  is s.t.,

$$\begin{aligned} \int_F \gamma_{t,F \underline{\mathbf{v}}_F}^k \cdot (\mathbf{rot}_F r + \mathbf{w}) \\ = \int_F \mathbf{C}_{F \underline{\mathbf{v}}_F}^k r + \sum_{E \in \mathcal{E}_F} \omega_{FE} \int_E \mathbf{v}_E r + \int_F \mathbf{v}_F \cdot \mathbf{w} \\ \forall (r, \mathbf{w}) \in \mathcal{P}^{0,k+1}(F) \times \mathcal{R}^{c,k}(F) \end{aligned}$$

- We have the following **polynomial consistency**:

$$C_F^k(\underline{\mathbf{I}}_{\text{rot},F}^k \mathbf{v}) = \mathbf{rot}_F \mathbf{v} \quad \forall \mathbf{v} \in \mathcal{N}^{k+1}(F) \quad \text{and} \quad \gamma_{t,F}^k(\underline{\mathbf{I}}_{\text{rot},F}^k \mathbf{v}) = \mathbf{v} \quad \forall \mathbf{v} \in \mathcal{P}^k(F)$$



# Two-dimensional DDR complex

Space	$V$ (vertex)	$E$ (edge)	$F$ (face)
$\underline{X}_{\text{grad},F}^k$	$\mathbb{R}$	$\mathcal{P}^{k-1}(E)$	$\mathcal{P}^{k-1}(F)$
$\underline{X}_{\text{curl},F}^k$		$\mathcal{P}^k(E)$	$\mathcal{RT}^k(F)$
$\mathcal{P}^k(F)$			$\mathcal{P}^k(F)$

- Define the **discrete gradient**

$$\underline{G}_F^k q_F := (\pi_{\mathcal{RT},F}^k \mathbf{G}_F^k q_F, (G_E^k q_F)_{E \in \mathcal{E}_F})$$

- The **two-dimensional DDR complex** reads

$$\underline{X}_{\text{grad},F}^k \xrightarrow{\underline{G}_F^k} \underline{X}_{\text{curl},F}^k \xrightarrow{C_F^k} \mathcal{P}^k(F) \longrightarrow \{0\}$$

- If  $F$  is simply connected, this complex is **exact**

# A glance at the general case I

- $\underline{X}_h^{k,r}$  spanned by vectors of polynomial components
- Recursive and hierarchical construction on  $d$ -cells,  $d = r + 1, \dots, n$ , of
  - A **discrete exterior derivative**

$$d_f^{k,r} : \underline{X}_f^{k,r} \rightarrow \mathcal{P}^k \Lambda^{r+1}(f)$$

- Based on it, an associated **discrete potential** ( $\simeq k$ -form inside  $f$ )

$$P_f^{k,r} : \underline{X}_f^{k,r} \rightarrow \mathcal{P}^k \Lambda^r(f)$$

- Reconstructions mimic the **Stokes formula**:  $\forall (\omega, \mu) \in \Lambda^\ell(f) \times \Lambda^{n-\ell-1}(f)$ ,

$$\int_f d^\ell \omega \wedge \mu = (-1)^{\ell+1} \int_f \omega \wedge d^{n-\ell-1} \mu + \int_{\partial f} \text{tr}_{\partial f} \omega \wedge \text{tr}_{\partial f} \mu$$

# A glance at the general case II

- For a polytopal domain  $\Omega \subset \mathbb{R}^n$  and a form degree  $r$ , the DDR space is

$$\underline{X}_h^{k,r} := \bigtimes_{d=r}^n \bigtimes_{f \in \Delta_d(\mathcal{T}_h)} \mathcal{P}^{-,k} \Lambda^{d-r}(f) \text{ with } \Delta_d(\mathcal{T}_h) := \{d\text{-faces of } \mathcal{T}_h\}$$

- We recursively define, for  $f \in \Delta_d(\mathcal{T}_h)$ ,  $d = r, \dots, n$ ,

- If  $r = d$ ,

$$\underline{P}_f^{k,d} \underline{\omega}_f := \star^{-1} \omega_f \in \mathcal{P}^k \Lambda^d(f)$$

- If  $r+1 \leq d \leq n$ , we first let, for all  $\underline{\omega}_f \in \underline{X}_f^{k,r}$  and all  $\mu \in \mathcal{P}^k \Lambda^{d-r-1}(f)$ ,

$$\int_f \underline{d}_f^{k,r} \underline{\omega}_f \wedge \mu = (-1)^{r+1} \int_f \star^{-1} \omega_f \wedge d\mu + \int_{\partial f} \underline{P}_{\partial f}^{k,r} \underline{\omega}_{\partial f} \wedge \text{tr}_{\partial f} \mu$$

then, for all  $(\mu, \nu) \in \kappa \mathcal{P}^{k,d-r}(f) \times \kappa \mathcal{P}^{k-1,d-r+1}(f)$ ,

$$\begin{aligned} (-1)^{r+1} \int_f \underline{P}_f^{k,r} \underline{\omega}_f \wedge (d\mu + \nu) &= \int_f \underline{d}_f^{k,r} \underline{\omega}_f \wedge \mu \\ &\quad - \int_{\partial f} \underline{P}_{\partial f}^{k,r} \underline{\omega}_{\partial f} \wedge \text{tr}_{\partial f} \mu + (-1)^{r+1} \int_f \star^{-1} \omega_f \wedge \nu \end{aligned}$$

# A glance at the general case III

- The following **polynomial consistency properties** hold:

$$\begin{aligned}P_f^{k,r} \underline{I}_f^{k,r} \omega &= \omega \quad \forall \omega \in \mathcal{P}^k \Lambda^r(f), \\ \underline{d}_f^{k,r} \underline{I}_f^{k,r} \omega &= \underline{d} \omega \quad \forall \omega \in \mathcal{P}^{-,k+1} \Lambda^r(f)\end{aligned}$$

- Setting

$$\underline{d}_h^{k,r} \underline{\omega}_h := (\pi_f^{-,k,d-r-1}(\star \underline{d}_f^{k,r} \underline{\omega}_f))_{f \in \Delta_d(\mathcal{T}_h), d \in [k+1, n]},$$

the **global DDR complex of differential forms** reads

$$\underline{X}_h^{k,0} \xrightarrow{\underline{d}_h^{k,0}} \underline{X}_h^{k,1} \longrightarrow \cdots \longrightarrow \underline{X}_h^{k,n-1} \xrightarrow{\underline{d}_h^{k,n-1}} \underline{X}_h^{k,n} \longrightarrow \{0\}$$

# A glance at the general case IV

For  $n = 3$ , vector proxies yield the DDR complex of [DP and Droniou, 2023a]:

$$\underline{X}_{\text{grad},T}^k \xrightarrow{\underline{G}_T^k} \underline{X}_{\text{curl},T}^k \xrightarrow{\underline{C}_T^k} \underline{X}_{\text{div},T}^k \xrightarrow{D_T^k} \mathcal{P}^k(T) \longrightarrow \{0\}$$

Space	$V$	$E$	$F$	$T$ (element)
$\underline{X}_T^{k,0} \cong \underline{X}_{\text{grad},T}^k$	$\mathbb{R}$	$\mathcal{P}^{k-1}(E)$	$\mathcal{P}^{k-1}(F)$	$\mathcal{P}^{k-1}(T)$
$\underline{X}_T^{k,1} \cong \underline{X}_{\text{curl},T}^k$		$\mathcal{P}^k(E)$	$\mathcal{RT}^k(F)$	$\mathcal{RT}^k(T)$
$\underline{X}_T^{k,2} \cong \underline{X}_{\text{div},T}^k$			$\mathcal{P}^k(F)$	$\mathcal{N}^k(T)$
$\underline{X}_T^{k,3} \cong \mathcal{P}^k(T)$				$\mathcal{P}^k(T)$

# Commutation with the interpolators

## Lemma (Local commutation properties)

The following diagrams commute:

$$\begin{array}{ccccccccc}
 C^\infty(\bar{T}) & \xrightarrow{\text{grad}} & C^\infty(\bar{T})^3 & \xrightarrow{\text{curl}} & C^\infty(\bar{T})^3 & \xrightarrow{\text{div}} & C^\infty(\bar{T}) & \longrightarrow & \{0\} \\
 \downarrow \underline{I}_{\text{grad},T}^k & & \downarrow \underline{I}_{\text{curl},T}^k & & \downarrow \underline{I}_{\text{div},T}^k & & \downarrow i_T & & \\
 \underline{X}_{\text{grad},T}^k & \xrightarrow{\underline{G}_T^k} & \underline{X}_{\text{curl},T}^k & \xrightarrow{\underline{C}_T^k} & \underline{X}_{\text{div},T}^k & \xrightarrow{D_T^k} & \mathcal{P}^k(T) & \longrightarrow & \{0\}
 \end{array}$$

- Crucial for both algebraic and analytical properties
- Compatibility of projections with **Helmholtz–Hodge decompositions**  
 $\implies$  robustness of DDR schemes with respect to the physics, e.g.:
  - Stokes [Beirão da Veiga, Dassi, DP, Droniou, 2022]
  - Navier–Stokes [DP, Droniou, Qian, 2023]
  - Reissner–Mindlin [DP and Droniou, 2023b]

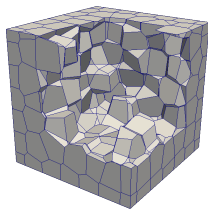
# Local discrete $L^2$ -products

- Based on the element potentials, we construct **local discrete  $L^2$ -products**

$$(x_T, y_T)_{\bullet, T} = \underbrace{\int_T P_{\bullet, T} x_T \cdot P_{\bullet, T} y_T}_{\text{consistency}} + \underbrace{s_{\bullet, T}(x_T, y_T)}_{\text{stability}} \quad \forall \bullet \in \{\text{grad, curl, div}\}$$

- The  $L^2$ -products are built to be **polynomially consistent**

# Global DDR complex



$$\underline{X}_{\text{grad},h}^k \xrightarrow{\underline{G}_h^k} \underline{X}_{\text{curl},h}^k \xrightarrow{\underline{C}_h^k} \underline{X}_{\text{div},h}^k \xrightarrow{D_h^k} \mathcal{P}^k(\mathcal{T}_h) \longrightarrow \{0\}$$

- **Global DDR spaces** on a mesh  $\mathcal{T}_h$  are defined gluing boundary components
- **Global operators** are obtained collecting local components
- **Global  $L^2$ -products  $(\cdot, \cdot)_{\bullet,h}$**  are obtained assembling element-wise



# Cohomology of the global three-dimensional DDR complex

$$\underline{X}_{\text{grad},h}^k \xrightarrow{\underline{G}_h^k} \underline{X}_{\text{curl},h}^k \xrightarrow{\underline{C}_h^k} \underline{X}_{\text{div},h}^k \xrightarrow{D_h^k} \mathcal{P}^k(\mathcal{T}_h) \longrightarrow \{0\}$$

Theorem (Cohomology of the 3D DDR complex [DP, Droniou, Pitassi, 2023])

For any  $k \geq 0$ , the DDR sequence forms a complex whose *cohomology spaces are isomorphic to those of the continuous de Rham complex*. In particular, if  $\Omega$  has a trivial topology (i.e.,  $b_1 = b_2 = 0$ ), the DDR complex is *exact*, i.e.,

$$\text{Im } \underline{G}_h^k = \text{Ker } \underline{C}_h^k, \quad \text{Im } \underline{C}_h^k = \text{Ker } D_h^k, \quad \text{Im } D_h^k = \mathcal{P}^k(\mathcal{T}_h).$$

Remark (Extension to PEC [Bonaldi, DP, Droniou, Hu, 2023])

The above result extends to the de Rham complex of differential forms.

# Outline

- 1 Three model problems and their well-posedness
- 2 Discrete de Rham (DDR) complexes
- 3 Application to magnetostatics**

# Uniform discrete Poincaré inequality for the curl

Theorem (Poincaré inequality for the curl [DP and Hanot, 2024])

For all  $\underline{\mathbf{v}}_h \in \underline{\mathbf{X}}_{\text{curl},h}^k$ , it holds

$$\inf_{\underline{\mathbf{z}}_h \in \text{Ker } \underline{\mathbf{C}}_h^k} \|\underline{\mathbf{v}}_h - \underline{\mathbf{z}}_h\|_{\text{curl},h} \lesssim \|\underline{\mathbf{C}}_h^k \underline{\mathbf{v}}_h\|_{\text{div},h},$$

with hidden constant only depending on  $\Omega$ , mesh regularity, and  $k$ .

**This results holds for domains of general topology!**

# Adjoint consistency of the discrete curl

**Adjoint consistency** measures the failure to satisfy a global IBP. For the curl,

$$\int_{\Omega} \mathbf{w} \cdot \mathbf{curl} \, \mathbf{v} - \int_{\Omega} \mathbf{curl} \, \mathbf{w} \cdot \mathbf{v} = 0 \text{ if } \mathbf{w} \times \mathbf{n} = \mathbf{0} \text{ on } \partial\Omega$$

## Theorem (Adjoint consistency for the curl)

Let  $\mathcal{E}_{\text{curl},h} : (C^0(\overline{\Omega})^3 \cap \mathbf{H}_0(\mathbf{curl}; \Omega)) \times \underline{\mathbf{X}}_{\text{curl},h}^k \rightarrow \mathbb{R}$  be s.t.

$$\mathcal{E}_{\text{curl},h}(\mathbf{w}, \underline{\mathbf{v}}_h) := (\underline{\mathbf{I}}_{\text{div},h}^k \mathbf{w}, \underline{\mathbf{C}}_h^k \underline{\mathbf{v}}_h)_{\text{div},h} - \int_{\Omega} \mathbf{curl} \, \mathbf{w} \cdot \mathbf{P}_{\text{curl},h}^k \underline{\mathbf{v}}_h.$$

Then, for all  $\mathbf{w} \in C^0(\overline{\Omega})^3 \cap \mathbf{H}_0(\mathbf{curl}; \Omega)$  s.t.  $\mathbf{w} \in H^{k+2}(\mathcal{T}_h)^3$ :  $\forall \underline{\mathbf{v}}_h \in \underline{\mathbf{X}}_{\text{curl},h}^k$ ,

$$|\mathcal{E}_{\text{curl},h}(\mathbf{w}, \underline{\mathbf{v}}_h)| \lesssim h^{k+1} \left( \|\underline{\mathbf{v}}_h\|_{\text{curl},h} + \|\underline{\mathbf{C}}_h^k \underline{\mathbf{v}}_h\|_{\text{div},h} \right).$$

# DDR scheme

- Assume  $b_2 = 0$ . We seek  $(\mathbf{H}, \mathbf{A}) \in \mathbf{H}(\text{curl}; \Omega) \times \mathbf{H}(\text{div}; \Omega)$  s.t.

$$\begin{aligned} \int_{\Omega} \mu \mathbf{H} \cdot \boldsymbol{\tau} - \int_{\Omega} \mathbf{A} \cdot \text{curl } \boldsymbol{\tau} &= 0 & \forall \boldsymbol{\tau} \in \mathbf{H}(\text{curl}; \Omega), \\ \int_{\Omega} \text{curl } \mathbf{H} \cdot \mathbf{v} + \int_{\Omega} \text{div } \mathbf{A} \text{ div } \mathbf{v} &= \int_{\Omega} \mathbf{J} \cdot \mathbf{v} & \forall \mathbf{v} \in \mathbf{H}(\text{div}; \Omega) \end{aligned}$$

- With obvious substitutions: Find  $(\underline{\mathbf{H}}_h, \underline{\mathbf{A}}_h) \in \underline{\mathbf{X}}_{\text{curl},h}^k \times \underline{\mathbf{X}}_{\text{div},h}^k$  s.t.

$$\begin{aligned} (\mu \underline{\mathbf{H}}_h, \underline{\boldsymbol{\tau}}_h)_{\text{curl},h} - (\underline{\mathbf{A}}_h, \underline{\mathbf{C}}_h^k \underline{\boldsymbol{\tau}}_h)_{\text{div},h} &= 0 & \forall \underline{\boldsymbol{\tau}}_h \in \underline{\mathbf{X}}_{\text{curl},h}^k, \\ (\underline{\mathbf{C}}_h^k \underline{\mathbf{H}}_h, \underline{\mathbf{v}}_h)_{\text{div},h} + \int_{\Omega} D_h^k \underline{\mathbf{A}}_h D_h^k \underline{\mathbf{v}}_h &= l_h(\underline{\mathbf{v}}_h) & \forall \underline{\mathbf{v}}_h \in \underline{\mathbf{X}}_{\text{div},h}^k \end{aligned}$$

- If  $b_2 \neq 0$ , we need to add **orthogonality of  $\underline{\mathbf{A}}_h$  to harmonic forms**

$$\underline{\mathfrak{H}}_{\text{div},h}^k := \left\{ \underline{\mathbf{w}}_h \in \underline{\mathbf{X}}_{\text{div},h}^k : D_h^k \underline{\mathbf{w}}_h = 0 \text{ and } (\underline{\mathbf{w}}_h, \underline{\mathbf{C}}_h^k \underline{\mathbf{v}}_h)_{\text{div},h} = 0 \forall \underline{\mathbf{v}}_h \in \underline{\mathbf{X}}_{\text{curl},h}^k \right\}$$

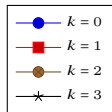
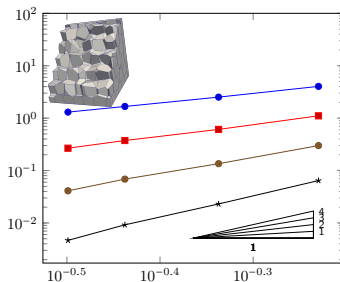
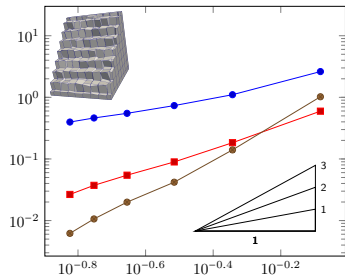
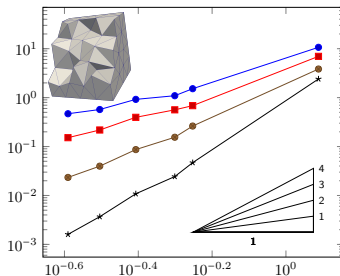
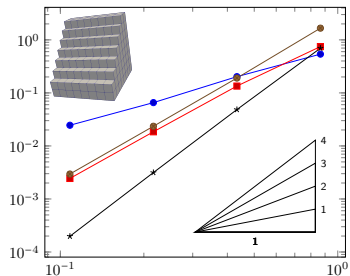
- **Inf-sup stability** is proved as in the continuous case for the norm

$$\|(\underline{\boldsymbol{\tau}}_h, \underline{\boldsymbol{v}}_h)\|_h := \left( \|\underline{\boldsymbol{\tau}}_h\|_{\text{curl},h}^2 + \|\underline{\boldsymbol{C}}_h^k \underline{\boldsymbol{\tau}}_h\|_{\text{div},h}^2 + \|\underline{\boldsymbol{v}}_h\|_{\text{div},h}^2 + \|D_h^k \underline{\boldsymbol{v}}_h\|_{L^2(\Omega)}^2 \right)^{\frac{1}{2}}$$

- Crucial points are the **isomorphism in cohomology** and **Poincaré inequality**
- Assuming  $\boldsymbol{H} \in C^0(\overline{\Omega})^3 \cap H^{k+2}(\mathcal{T}_h)^3$  and  $\boldsymbol{A} \in C^0(\overline{\Omega})^3 \times H^{k+2}(\mathcal{T}_h)^3$ , it holds

$$\|(\underline{\boldsymbol{H}}_h - \underline{\boldsymbol{I}}_{\text{curl},h}^k \boldsymbol{H}, \underline{\boldsymbol{A}}_h - \underline{\boldsymbol{I}}_{\text{div},h}^k \boldsymbol{A})\|_h \lesssim h^{k+1}$$

# Numerical examples (energy error vs. meshsize)



# Conclusions and perspectives

- Fully discrete approach for PDEs relating to the de Rham complex
- Key features: support of general polyhedral meshes and high-order
- Novel computational strategies made possible
- Natural extensions to differential forms
- Unified proof of analytical properties using differential forms
- Development of novel complexes (e.g., elasticity, Hessian, . . . )
- Applications (possibly beyond continuum mechanics)



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