

New generation numerical methods for partial differential equations



Daniele Di Pietro, 19 June 2024



UNIVERSITÉ DE
MONTPELLIER



ERC Synergy Grants

- **Largest** and **most competitive** grants awarded by the ERC
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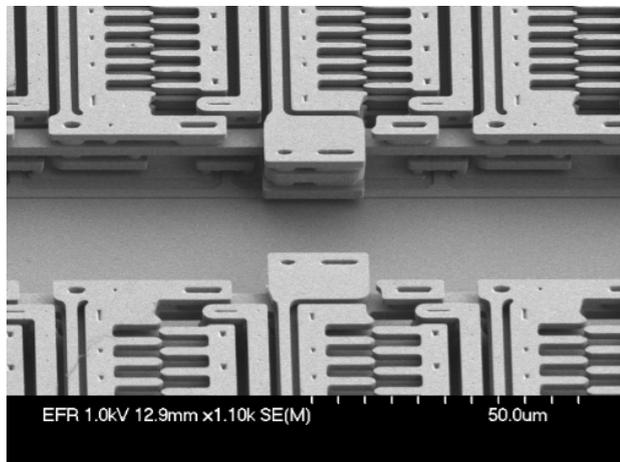
The NEMESIS project

- Principal Investigators (all researchers in Numerical Analysis!):
 - **Daniele Di Pietro**, Université de Montpellier, IMAG, corr. PI
 - **Paola Antonietti**, Politecnico di Milano, MOX
 - **Lourenço Beirão da Veiga**, Università di Milano Bicocca
 - **Jérôme Droniou**, CNRS, IMAG
- **7.8M€** (**4.4M€** at IMAG), **4** research clusters, **8** work packages
- **>70 + >15 man-year** of non-permanent + permanent researchers



From physical problems to numerical simulations

I. Physical problem



II. Mathematical model

$$A(u) = f$$



$$\begin{aligned} \mu H - \text{curl } A &= 0 && \text{in } \Omega, \\ \text{curl } H &= J && \text{in } \Omega, \\ \text{div } A &= 0 && \text{in } \Omega, \\ A \times n &= 0 && \text{on } \partial\Omega \end{aligned}$$

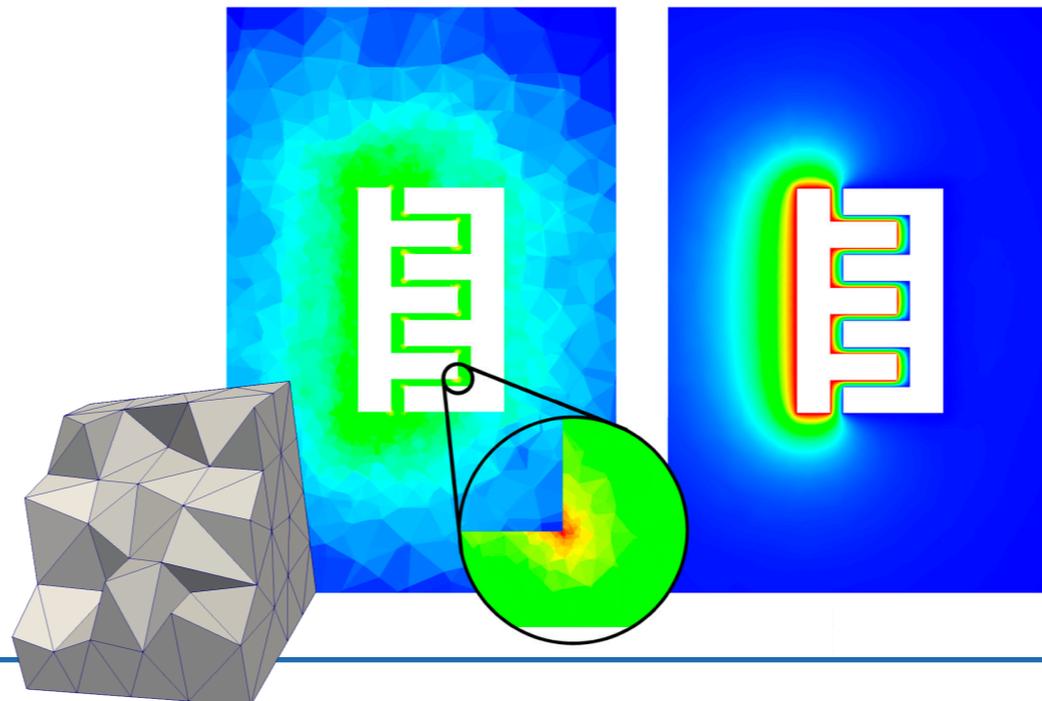
III. Numerical approximation

$$A_h(u_h) = f_h$$



IV. Computer simulation

$$A_h u_h = f_h$$



 NEMESIS



The challenges for next generation simulators

01.

↓ Incomplete differential operators in Hilbert complexes

02.

↓ Efficient solution of large, indefinite algebraic systems

03.

↓ Nonlinear, hybrid-dimensional physics

04.

↓ Handling all of the above at once

These challenges correspond to the **research clusters** of NEMESIS



Stability, consistency, and convergence

$$A_h(u_h) = f_h$$

- $\frac{1}{h}$ measures the **effort required to solve the discrete problem**
- Our ultimate goal is to have **convergent schemes**, for which

$$u_h \rightarrow u \text{ as } h \rightarrow 0$$

- **Stability:** Small variations of f_h induce small variations of u_h
- **Consistency:** $f_h - A_h(I_h u) \rightarrow 0$ as $h \rightarrow 0$
- For linear problems, we have the **Lax principle***:

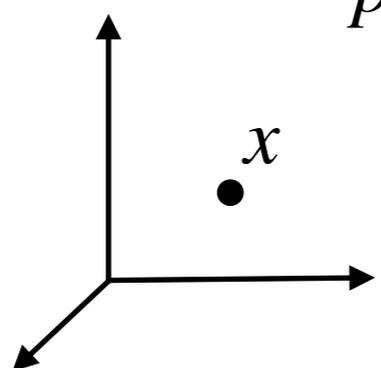
$$\text{Stability} \implies (\text{Consistency} \iff \text{Convergence})$$

* See, e.g., [Lax & Richtmyer, 1956]



Physical quantities have different nature

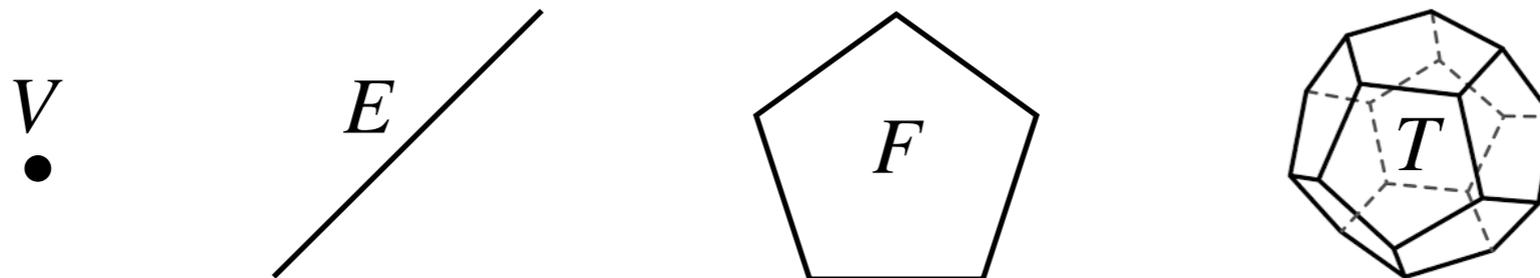
- **Scalar:** potential (pressure p) or density (energy \mathcal{E})
- **Vector:** circulation (magnetic field H) or flux (heat flux Φ)
- **Tensor:** (deformation ε , etc.)


$$p(x)$$
$$H(x) = \begin{pmatrix} H_1 \\ H_2 \\ H_3 \end{pmatrix}$$
$$\varepsilon(x) = \begin{pmatrix} \varepsilon_{11} & \varepsilon_{12} & \varepsilon_{13} \\ \varepsilon_{12} & \varepsilon_{22} & \varepsilon_{23} \\ \varepsilon_{13} & \varepsilon_{23} & \varepsilon_{33} \end{pmatrix}$$



How can we measure these quantities?

- Pressure p : evaluation at a point V
- Magnetic field H : free current in a wire E
- Heat flux Φ : normal flux through a surface F
- Energy \mathcal{E} : quantity in a volume T



Differential forms provide a unified approach to integration over curves, surfaces, solids, etc.

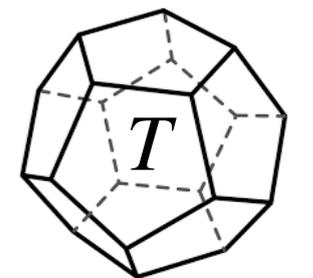
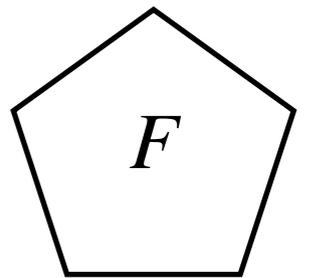
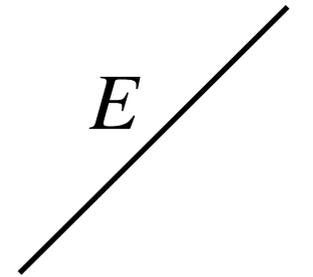


Vector calculus operators

$$p(x_{V_2}) - p(x_{V_1}) = \int_E \text{grad } p \cdot t_E d\ell \quad \text{grad } p = \begin{pmatrix} \partial_1 p \\ \partial_2 p \\ \partial_3 p \end{pmatrix}$$

$$-\oint_{\partial F} H \cdot t_{\partial F} d\ell = \int_F \text{curl } H \cdot n_F dS \quad \text{curl } H = \begin{pmatrix} \partial_2 H_3 - \partial_3 H_2 \\ \partial_3 H_1 - \partial_1 H_3 \\ \partial_1 H_2 - \partial_2 H_1 \end{pmatrix}$$

$$\oint_{\partial T} \Phi \cdot n_{\partial T} dS = \int_T \text{div } \Phi dV \quad \text{div } \Phi = \partial_1 \Phi_1 + \partial_2 \Phi_2 + \partial_3 \Phi_3$$



- curl and div are **incomplete differential operators**
- the notion of **exterior derivative** unifies grad, curl, and div



Betti numbers

- Let $\Omega \subset \mathbb{R}^3$ be a polyhedron with **Betti numbers** b_i
- $b_0 = 1$ (number of **connected components**) and $b_3 = 0$
- b_1 and b_2 account for the number of **tunnels** and **voids** in Ω



$$(b_0, b_1, b_2, b_3) = (1, 1, 0, 0)$$



$$(b_0, b_1, b_2, b_3) = (1, 0, 1, 0)$$



The de Rham cohomology

For a domain Ω of \mathbb{R}^3 , we can form the **de Rham complex**:

$$H^1(\Omega) \xrightarrow{\text{grad}} H(\text{curl}; \Omega) \xrightarrow{\text{curl}} H(\text{div}; \Omega) \xrightarrow{\text{div}} L^2(\Omega)$$

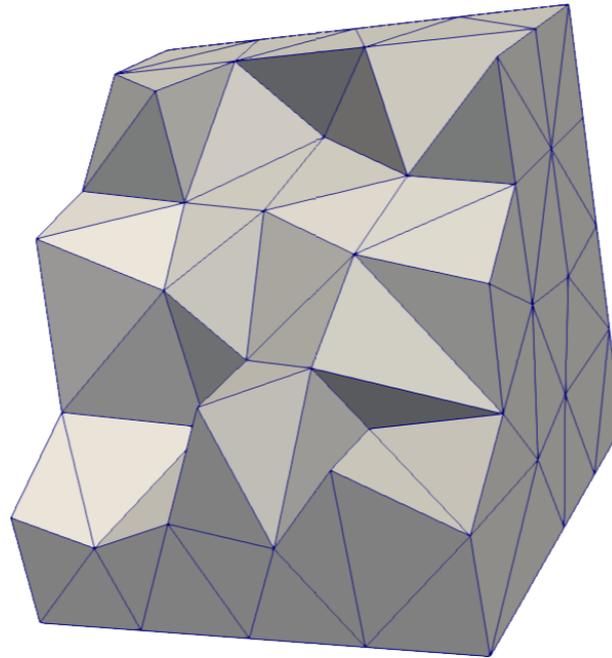
- Since **curl grad = 0** and **div curl = 0**, this is a **complex**
- Depending on Ω , we can strengthen these relations:
 - If $b_1 = 0$, **Ker curl = Im grad**
 - If $b_2 = 0$, **Im curl = Ker div**
- When $b_1 \neq 0$ or $b_2 \neq 0$, **de Rham's cohomology** characterizes

Ker curl/Im grad and Ker div/Im curl

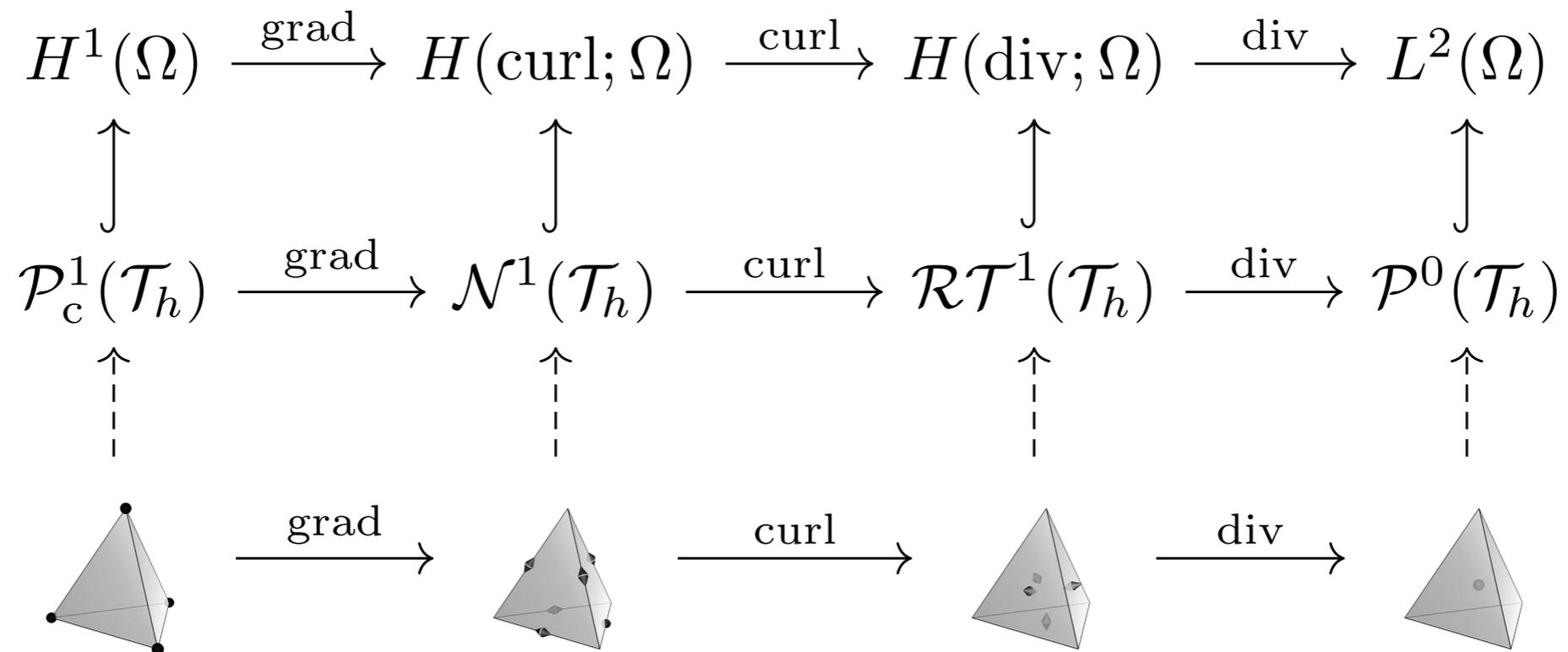
Discrete counterparts of these properties are key to **stability** when dealing with **incomplete differential operators!**



The Finite Element approach



A finite element mesh \mathcal{T}_h



[Raviart and Thomas, 1977], [Nédélec, 1980]

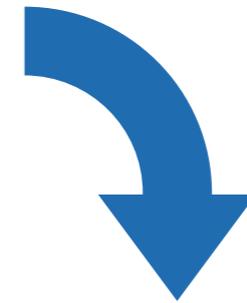


Finite Element discretization

Strong formulation

Find $H : \Omega \rightarrow \mathbb{R}^3$ and $A : \Omega \rightarrow \mathbb{R}^3$ s.t.

$$\begin{aligned}\mu H - \operatorname{curl} A &= 0 && \text{in } \Omega, \\ \operatorname{curl} H &= J && \text{in } \Omega, \\ \operatorname{div} A &= 0 && \text{in } \Omega, \\ A \times n &= 0 && \text{on } \partial\Omega\end{aligned}$$



Weak formulation

Find $H \in H(\operatorname{curl}; \Omega)$ and $A \in H(\operatorname{div}; \Omega)$ s.t.

$$\begin{aligned}\int_{\Omega} \mu H \cdot \tau - \int_{\Omega} A \cdot \operatorname{curl} \tau &= 0 && \forall \tau \in H(\operatorname{curl}; \Omega), \\ \int_{\Omega} \operatorname{curl} H \cdot v + \int_{\Omega} \operatorname{div} A \operatorname{div} v &= \int_{\Omega} f \cdot v && \forall v \in H(\operatorname{div}; \Omega)\end{aligned}$$

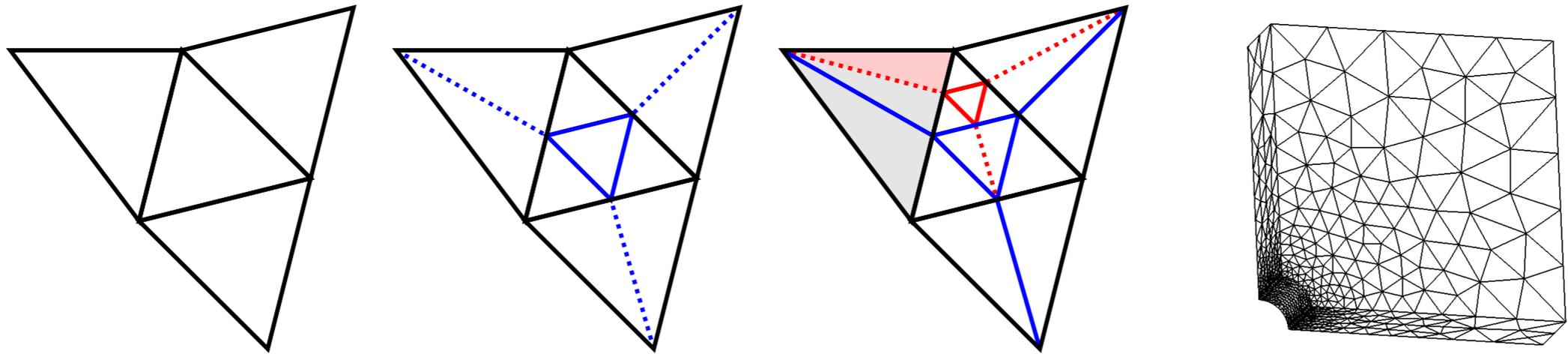
FE scheme

Find $H_h \in \mathcal{N}^1(\mathcal{T}_h)$ and $A_h \in \mathcal{RT}^1(\mathcal{T}_h)$ s.t.

$$\begin{aligned}\int_{\Omega} \mu H_h \cdot \tau_h - \int_{\Omega} A_h \cdot \operatorname{curl} \tau_h &= 0 && \forall \tau_h \in \mathcal{N}^1(\mathcal{T}_h), \\ \int_{\Omega} \operatorname{curl} H_h \cdot v_h + \int_{\Omega} \operatorname{div} A_h \operatorname{div} v_h &= \int_{\Omega} f \cdot v_h && \forall v_h \in \mathcal{RT}^1(\mathcal{T}_h)\end{aligned}$$



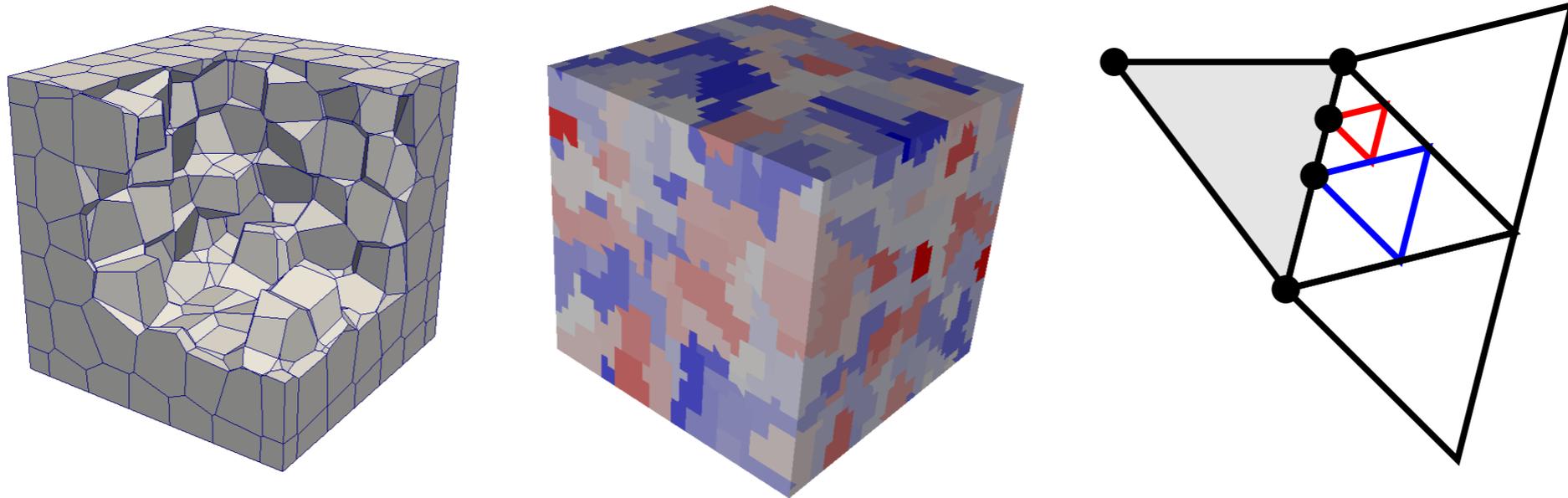
Limitations of Finite Elements



- Approach limited to **conforming meshes** with standard elements:
 - Local refinement requires to **trade mesh quality for size**
 - Complex geometries may require a **large number of elements**
 - The element shape cannot be **adapted to the solution**
- The extension to more advanced cases is not straightforward
- Solution of algebraic systems cannot benefit from **agglomeration**

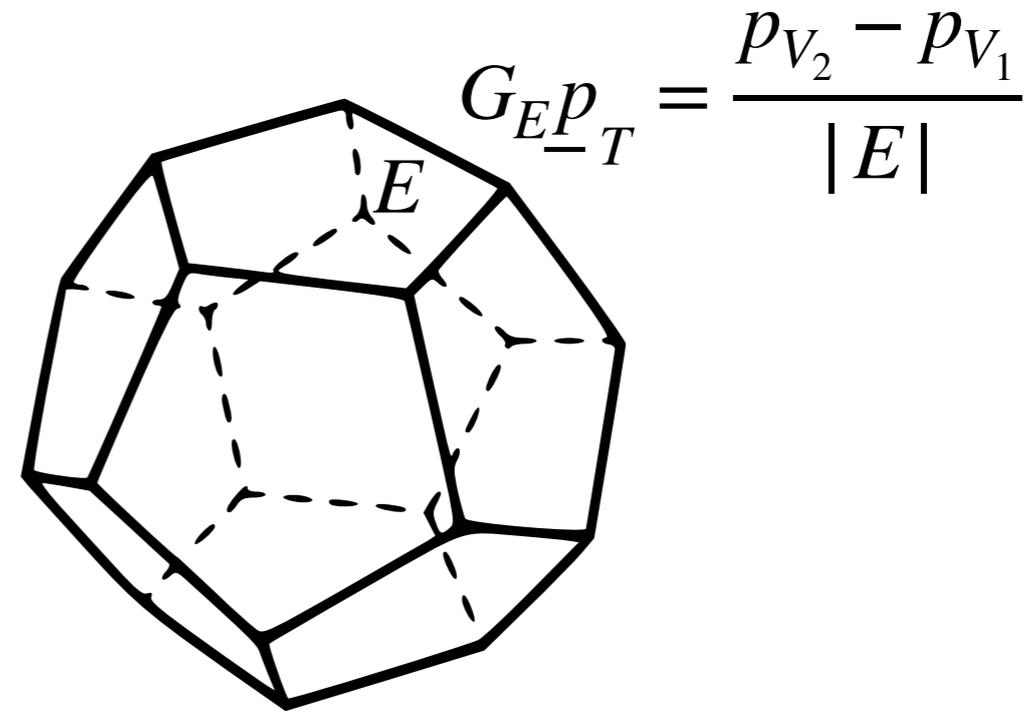
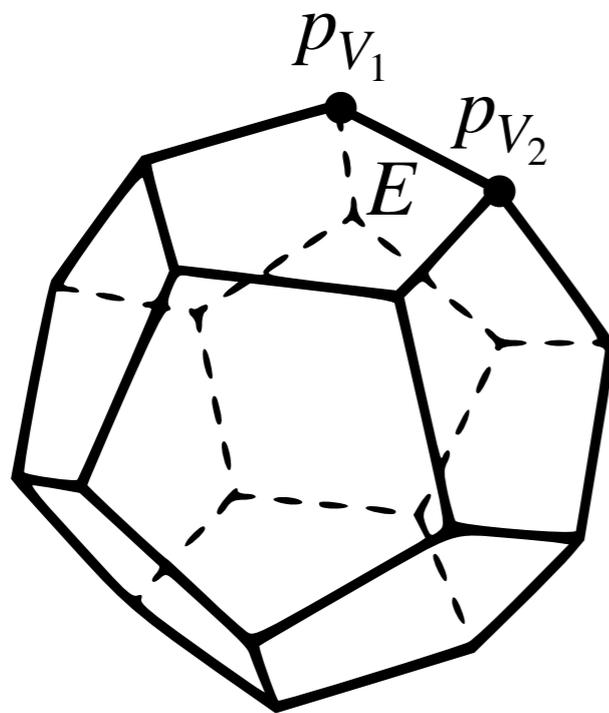


The fully discrete polyhedral approach



- **Key idea:** use discrete spaces *and* operators
- Support of general polyhedral meshes and high-order
- Several strategies to reduce the size of algebraic systems
- Agglomeration-based multi-grid solvers for algebraic systems

Discrete vector calculus operators

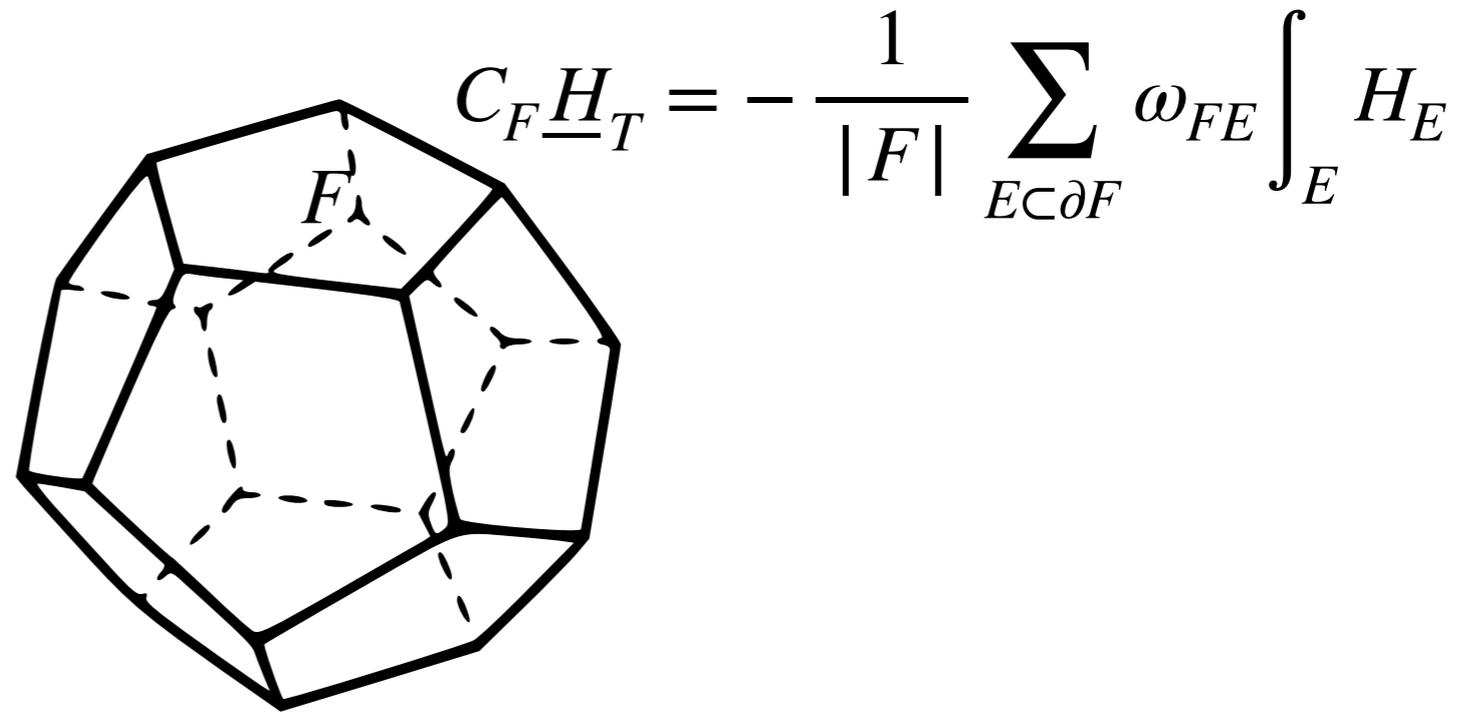
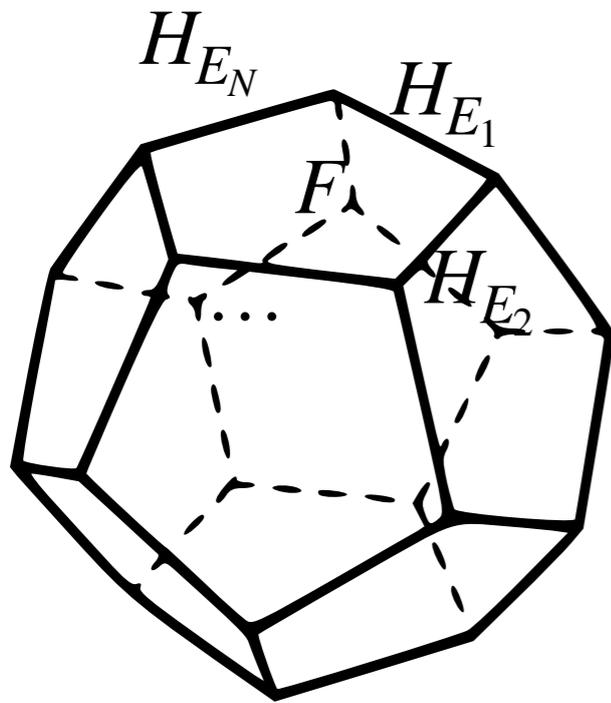


$$\underline{p}_T := (p_V)_{V \in \mathcal{V}_T} \in \mathbb{R}^{\mathcal{V}_T}$$

$$\underline{G}_T \underline{p}_T := (G_{E p_T})_{E \in \mathcal{E}_T} \in \mathbb{R}^{\mathcal{E}_T}$$

$$\underline{X}_{\text{grad}, T} := \mathbb{R}^{\mathcal{V}_T} \xrightarrow{\underline{G}_T} \underline{X}_{\text{curl}, T} := \mathbb{R}^{\mathcal{E}_T}$$

Discrete vector calculus operators



$$C_F \underline{H}_T = -\frac{1}{|F|} \sum_{E \subset \partial F} \omega_{FE} \int_E H_E$$

$$\underline{H}_T := (H_E)_{E \in \mathcal{E}_T} \in \mathbb{R}^{\mathcal{E}_T}$$

$$\underline{C}_T := (C_F \underline{H}_T)_{F \in \mathcal{F}_T} \in \mathbb{R}^{\mathcal{F}_T}$$

$$\underline{X}_{\text{grad},T} \xrightarrow{\underline{G}_T} \underline{X}_{\text{curl},T} \xrightarrow{\underline{C}_T} \underline{X}_{\text{div},T} := \mathbb{R}^{\mathcal{F}_T}$$



A lowest-order polyhedral scheme

$$\underline{X}_{\text{grad},h} \xrightarrow{\underline{G}_h} \underline{X}_{\text{curl},h} \xrightarrow{\underline{C}_h} \underline{X}_{\text{div},h} \xrightarrow{D_h} \mathbb{R}^{\mathcal{T}_h}$$

Weak formulation

Find $H \in H(\text{curl}; \Omega)$ and $A \in H(\text{div}; \Omega)$ s.t.

$$\int_{\Omega} \mu H \cdot \tau - \int_{\Omega} A \cdot \text{curl } \tau = 0 \quad \forall \tau \in H(\text{curl}; \Omega),$$

$$\int_{\Omega} \text{curl } H \cdot v + \int_{\Omega} \text{div } A \text{ div } v = \int_{\Omega} f \cdot v \quad \forall v \in H(\text{div}; \Omega)$$



Polyhedral scheme

Find $\underline{H}_h \in \underline{X}_{\text{curl},h}$ and $\underline{A}_h \in \underline{X}_{\text{div},h}$ s.t.

$$\mu(\underline{H}_h, \underline{\tau}_h)_{\text{curl},h} - (\underline{A}_h, \underline{C}_h \underline{\tau}_h)_{\text{div},h} = 0 \quad \forall \underline{\tau}_h \in \underline{X}_{\text{curl},h},$$

$$(\underline{C}_h \underline{H}_h, \underline{v}_h)_{\text{div},h} + \int_{\Omega} D_h \underline{A}_h D_h \underline{v}_h = \int_{\Omega} f \cdot P_{\text{div},h} \underline{v}_h \quad \forall \underline{v}_h \in \underline{X}_{\text{div},h}$$



Increasing the approximation order

- With the previous construction, one can hope for the error

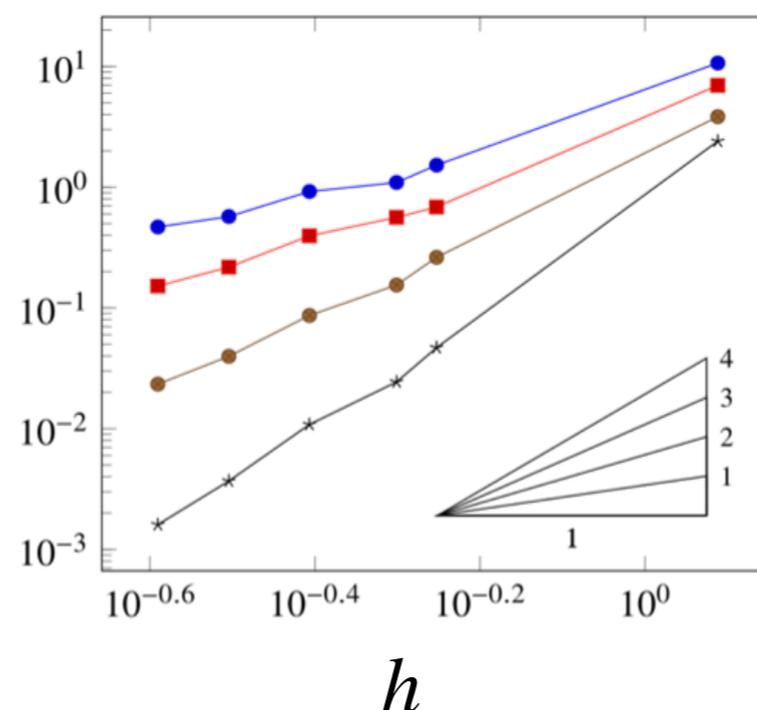
$$\|u - I_h u\|_h \lesssim h$$

- However, for u smooth, it would be desirable to have instead

$$\|u - I_h u\|_h \lesssim h^k \text{ with } k \geq 1$$

- This requires **high-order discrete de Rham complexes**

$$\|u_h - I_h u\|_h$$



Arbitrary order discrete de Rham complexes

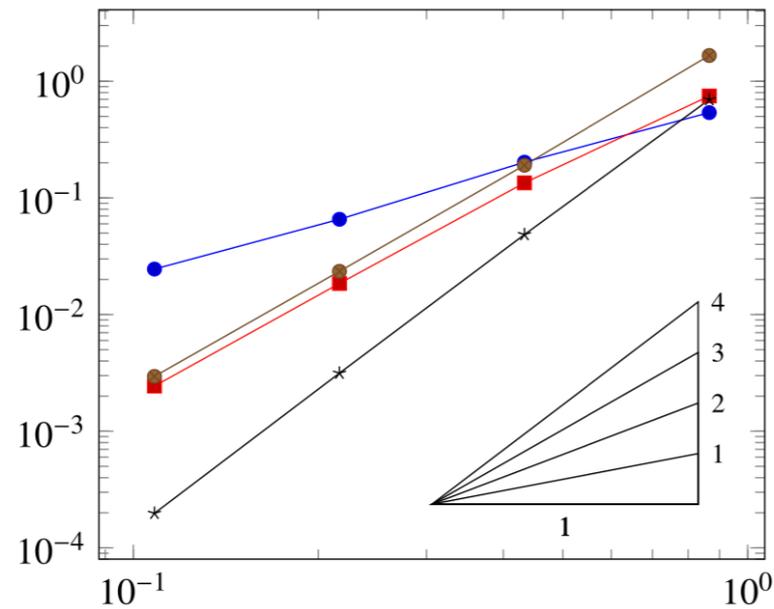
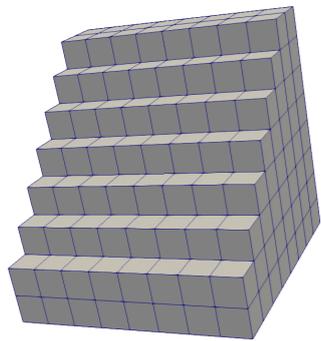
$$\underline{X}_{\text{grad},h}^k \xrightarrow{\underline{G}_h^k} \underline{X}_{\text{curl},h}^k \xrightarrow{\underline{C}_h^k} \underline{X}_{\text{div},h}^k \xrightarrow{D_h^k} \mathcal{P}^k(\mathcal{T}_h)$$

- First ideas based on FE [Beirão da Veiga et al., 2016–18]
- First DDR complex [DP, Droniou, and Rapetti, 2020]
- First complete set of analytical properties [DP and Droniou, 2023]
- Extension to differential forms [Bonaldi, DP, Droniou, Hu, 2023]
- **Key preliminary developments for the NEMESIS project!**

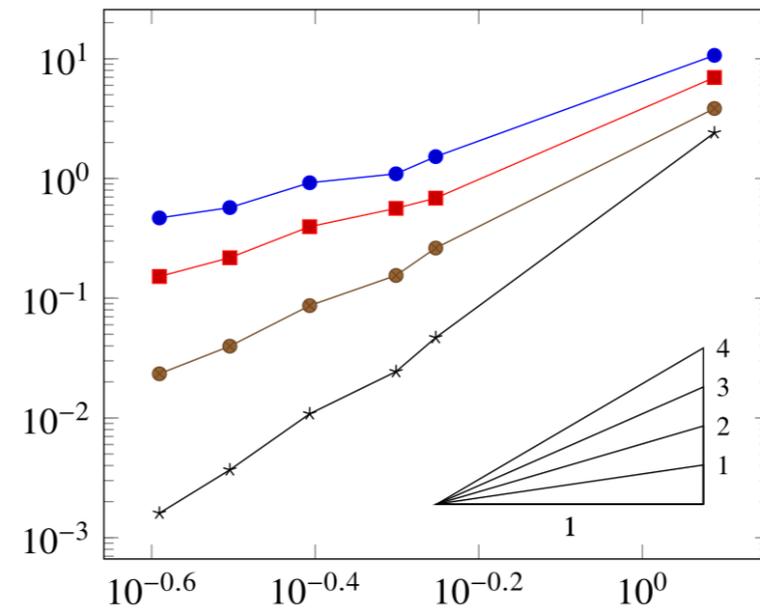


Convergence for the magnetostatics problem

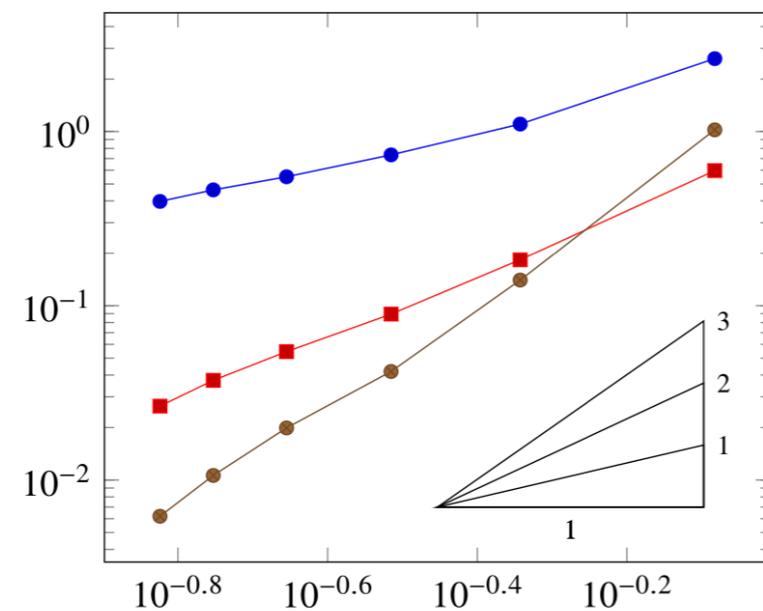
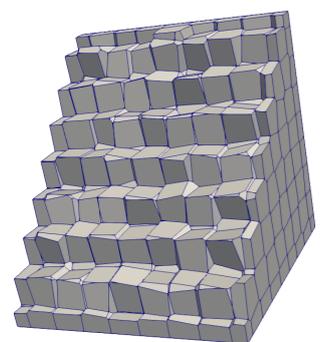
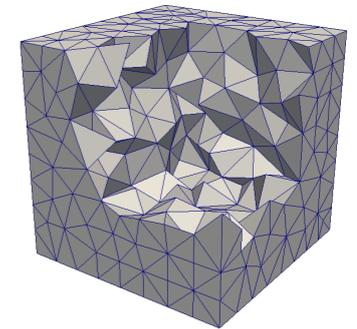
● $k = 0$
■ $k = 1$
● $k = 2$
* $k = 3$



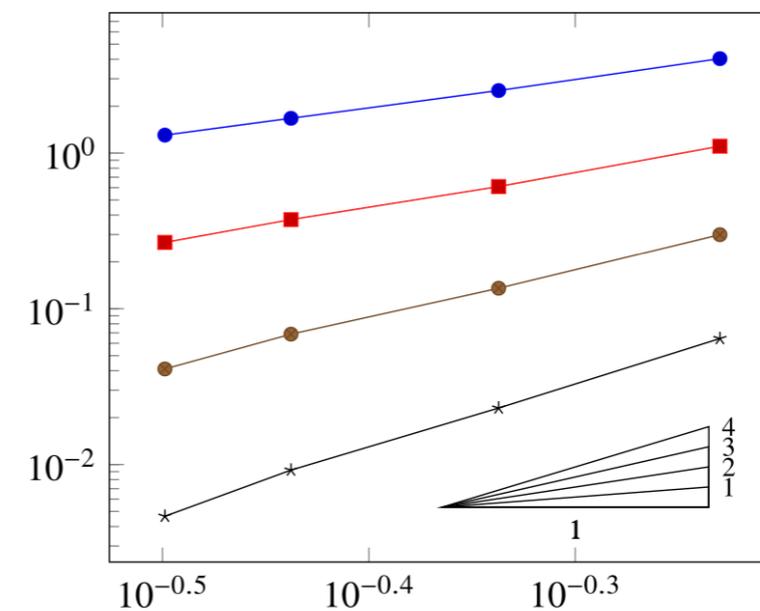
(a) "Cubic-Cells" mesh



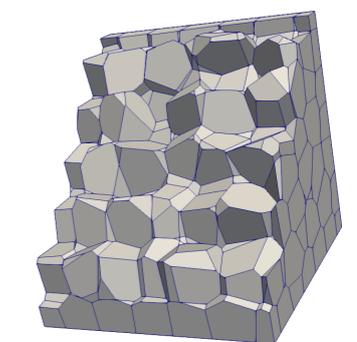
(b) "Tetgen-Cube-0" mesh



(c) "Voro-small-0" mesh



(d) "Voro-small-1" mesh



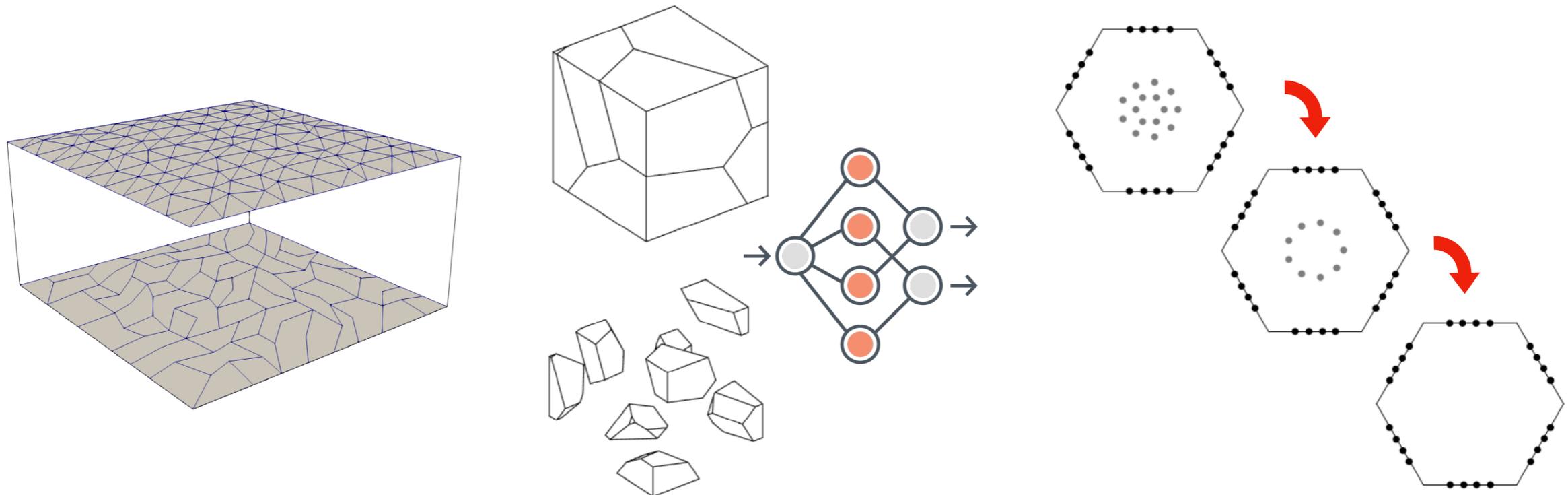
$\|(H_h - I_{\text{curl},h}^k H, A_h - I_{\text{div},h}^k A)\|_h$ vs. h for $k \in \{0,1,2,3\}$, [DP and Droniou, 2021]



Efficiency boost

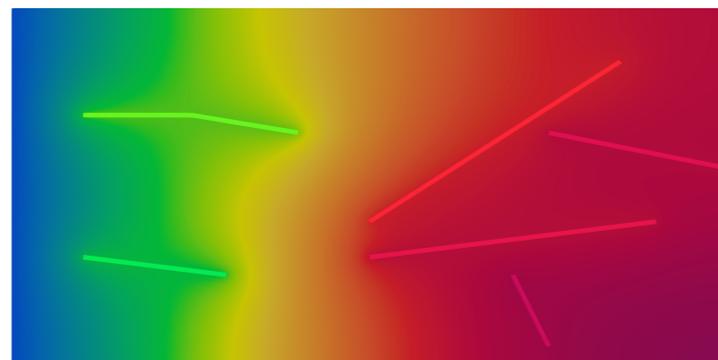
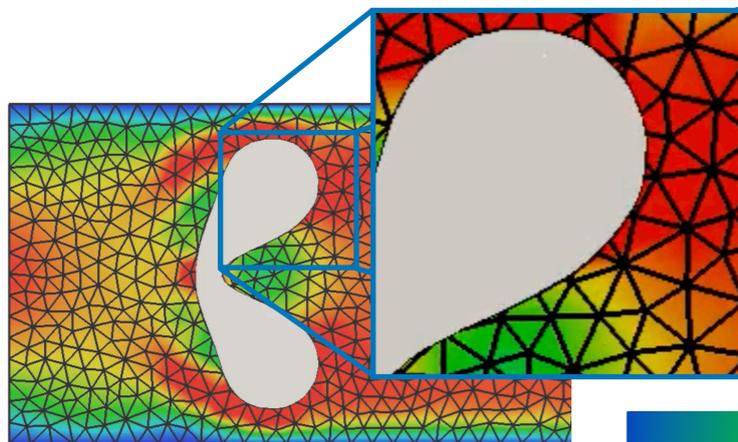
$$A_h u_h = f_h$$

- Agglomeration-based **multigrid solvers**
- AI-driven **mesh adaptation**
- Serendipity and static condensation to **reduce system size**



Hybrid-dimensional, coupled, nonlinear physics

- Rigorous mathematical tools for **highly non linear problems**
- Development of methods for **problems set on surfaces**
- Polyhedral meshes for **moving interface problems**



Discrete Rellich–Kondrachov on domains and manifolds:

$(D_h u_h)_{h \in \mathcal{H}}$ bounded $\implies (u_h)_{h \in \mathcal{H}}$ compact
with D_h discrete **curl**, **div**, **grad**_s, ...

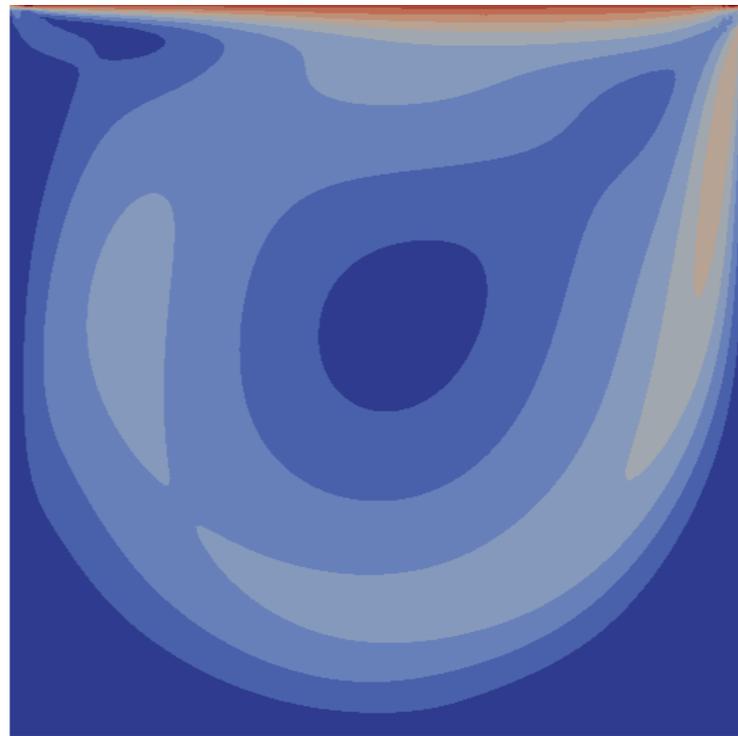


An example of highly nonlinear problem

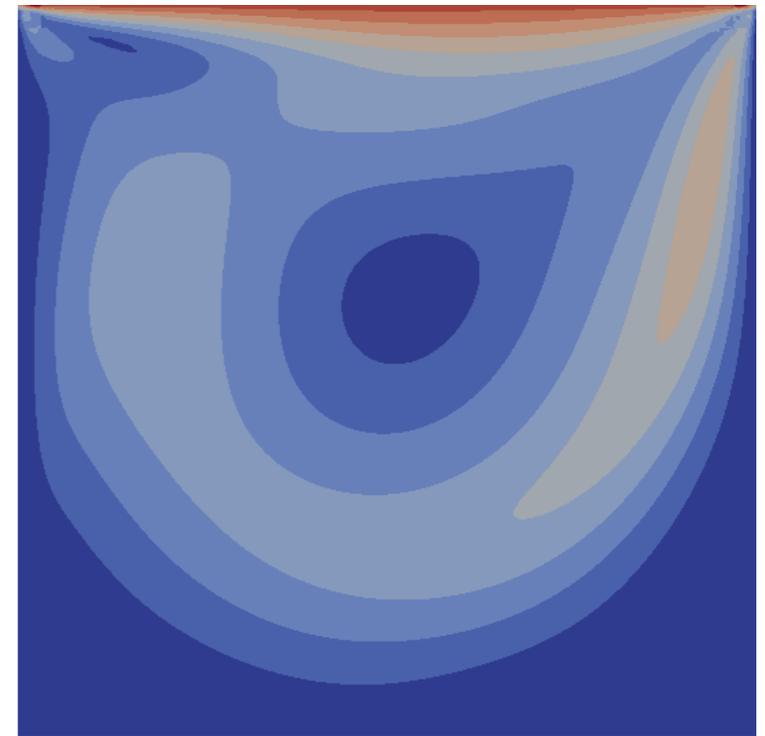
The non-Newtonian Navier–Stokes equations



$p = 1.5$



$p = 2$



$p = 2.5$

$$\begin{aligned} \partial_t u - \nabla \cdot \sigma(u) + (u \cdot \nabla)u + \nabla p &= f \\ \nabla \cdot u &= 0 \end{aligned} \quad \text{with } \sigma(u) = \nu |\varepsilon(u)|^{p-2} \varepsilon(u)$$

Example of mathematical challenges

- Development of **discrete elasticity complexes**
- Corresponding **discrete Poincaré and Sobolev inequalities**
- **Compactness results**
- ...

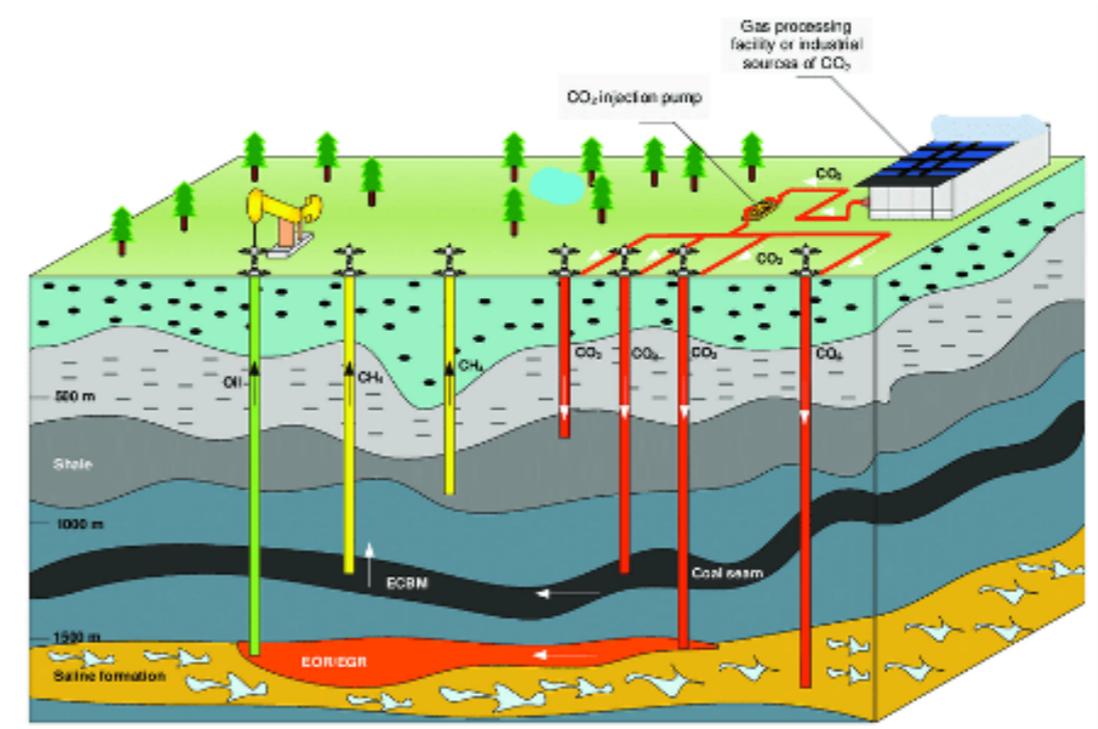
$$H^1(\Omega) \xrightarrow{\nabla_{\text{sym}}} H_{\mathbb{S}}(\text{Rot Rot}^{\top}; \Omega) \xrightarrow{\text{Rot Rot}^{\top}} H_{\mathbb{S}}(\text{Div}; \Omega) \xrightarrow{\text{Div}_{\mathbb{S}}} L^2(\Omega)$$



Proof-of-concept applications



Magnetohydrodynamics



Geological flows

These applications require to combine all the advances of the NEMESIS project!





New generation methods
for numerical simulations



Funded by
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European Research Council
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Thank you for your attention!

